

FIGURE 4.10: Polar radius t and mean radius β

is the geometric mean of all three axes (Fig. 4.10). (In a more familiar notation this is $R = \sqrt[3]{a^2b}$, the sphere being defined as having the same volume as the ellipsoid.) In view of the smallness of f , (4-210) reduces in the linear approximation to

$$\beta = t \left(1 + \frac{2}{3} f \right) , \quad (4-211)$$

$$t = \beta \left(1 - \frac{2}{3} f \right) . \quad (4-212)$$

Hence,

$$W'(t) = \frac{dW}{dt} = \frac{dW}{d\beta} \frac{d\beta}{dt} = \frac{dW}{d\beta} \left(1 + \frac{2}{3} f + \frac{2}{3} tf' \right) . \quad (4-213)$$

Using (4-209) with (4-211), this gives

$$W'(t) = -\frac{4\pi G}{3} Dt \left(1 + \frac{2}{3} f \right) \left(1 + \frac{2}{3} f + \frac{2}{3} tf' \right) + \frac{2}{3} \omega^2 t \quad (4-214)$$

(since $\omega^2 = O(f)$, we have been able simply to replace β by t in the last term).

Introducing the dimensionless quantity (4-66), in the present units

$$\mu = \frac{3}{4\pi G} \frac{\omega^2}{D} , \quad (4-215)$$

which is $O(f)$, we thus have to $O(f)$

$$g_P = -W'(t) = \frac{4\pi G}{3} Dt \left(1 + \frac{4}{3} f + \frac{2}{3} tf' - \frac{2}{3} \mu \right) . \quad (4-216)$$

Now

$$\frac{4\pi G \rho - 2\omega^2}{\frac{4\pi}{3} G t D} = \frac{3}{t} \frac{\delta}{D} - \frac{2}{t} \mu \quad (4-217)$$