

truncated at  $n = 2$ . This also shows that  $M$  is the total mass enclosed by the ellipsoid, which thus is seen to be equal to the mass of the auxiliary mean sphere of radius  $R$ .

This is quite normal since any ellipsoid  $E$  and its associated mean sphere  $S$  (of radius  $q$ ) enclose the same volume by the very definition (2-82), in view of (2-53) for  $n = 2$ : the mean deviation between  $E$  and  $S$  is zero. This holds for any ellipsoid of constant density,  $q < R$ , as well as for the boundary ellipsoid  $q = R$ , which we are considering in (2-92).

*Internal potential.* We shall use a similar artifice (Fig. 2.5) as for the sphere

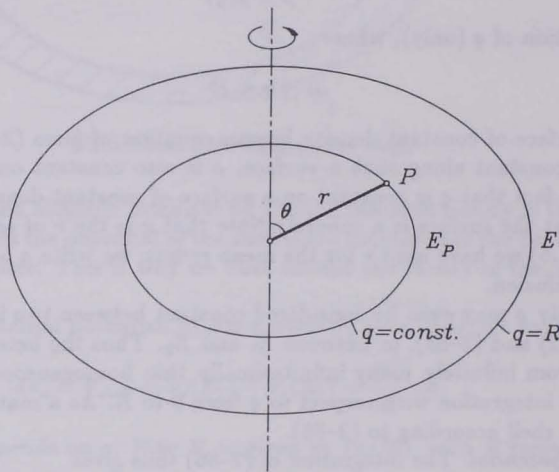


FIGURE 2.5: Illustrating the potential at an interior point  $P$

(Fig. 2.2), considering the ellipsoid (= ellipsoidal surface) of constant density  $E_P$  passing through the interior point  $P$  at which the potential  $V = V_i$  is to be computed. The ellipsoid  $E_P$  is characterized by its value  $q$  (the radius of the corresponding mean sphere); along  $E_P$ , the value of  $q$  is, of course, constant as we have already remarked. The equation of  $E_P$  is (2-82);  $r$  and  $\theta$  are shown in Fig. 2.5.

Again we shall build up the potential by summing (integrating) the contributions of the infinitesimal shells bounded by ellipsoids of constant density as shown in Fig. 2.4. These contributions are given by (2-86) and (2-87). Since  $q$  has been reserved for  $E_P$  (Fig. 2.5), we shall denote the integration variable by  $q'$ , similarly as we did for the sphere, cf. (2-47). For the interior of  $E_P$ , i.e. for  $q' < q$ , we take (2-86); for the shell between  $E_P$  and  $E$ , i.e. for  $q < q' < R$ , we take (2-87):  $P$  is external for the region inside  $E_P$  (being just on its external boundary  $E_P$ ) and internal for the shell. Thus we get